

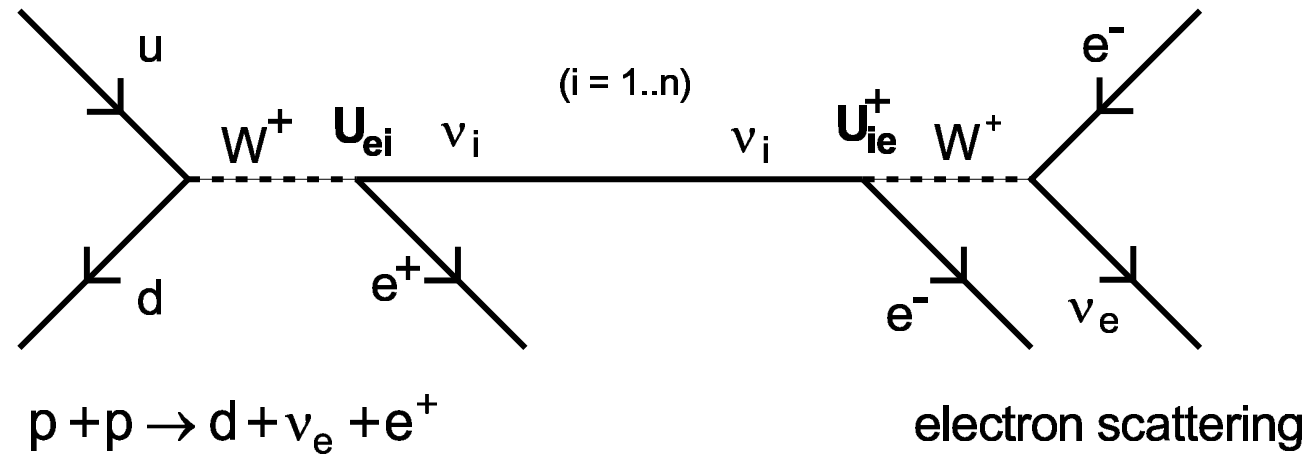
CP, T and CPT Violation at Future Long Baseline Experiments

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Neutrino Oscillations for $N = 2$

Production
Propagation
Detection



2 Flavours ν_e, ν_μ :

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

flavour states

mixing matrix · mass eigenstates

Probability:

$$P(\nu_\mu \rightarrow \nu_e) = |\langle \nu_\mu(t) | \nu_e(t=0) \rangle|^2 = \sin^2 2\theta \cdot \sin^2 \left(\frac{\Delta m^2 L}{4E_\nu} \right)$$

$N \geq 2$ with CP and Matter Effects

Precision: $N = 2$ description insufficient \Rightarrow modifications

- $2 \rightarrow 3$ neutrino framework \Rightarrow more parameters & CP effects
- MSW: parameter mapping in matter

$$\Rightarrow P(\nu_{e_l} \rightarrow \nu_{e_m}) = \underbrace{\delta_{lm} - 4 \sum_{i>j} \text{Re} J_{ij}^{e_l e_m} \sin^2 \Delta_{ij}}_{P_{CP}} - 2 \underbrace{\sum_{i>j} \text{Im} J_{ij}^{e_l e_m} \sin 2\Delta_{ij}}_{P_{CP}}$$

Shorthands: $J_{ij}^{e_l e_m} := U_{li} U_{lj}^* U_{mi}^* U_{mj}$ $\Delta_{ij} := \frac{\Delta m_{ij}^2 L}{4E}$

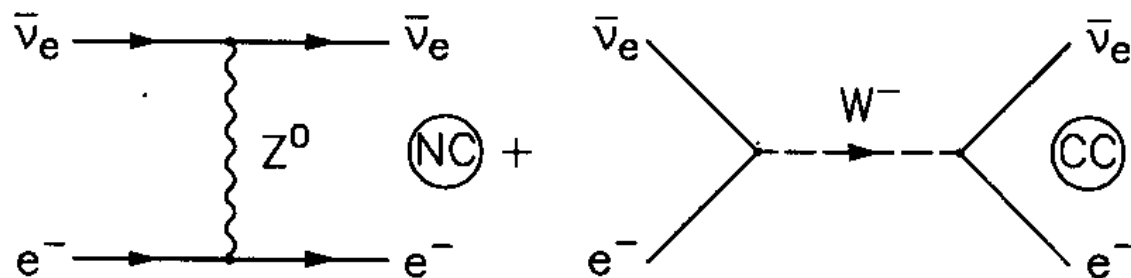
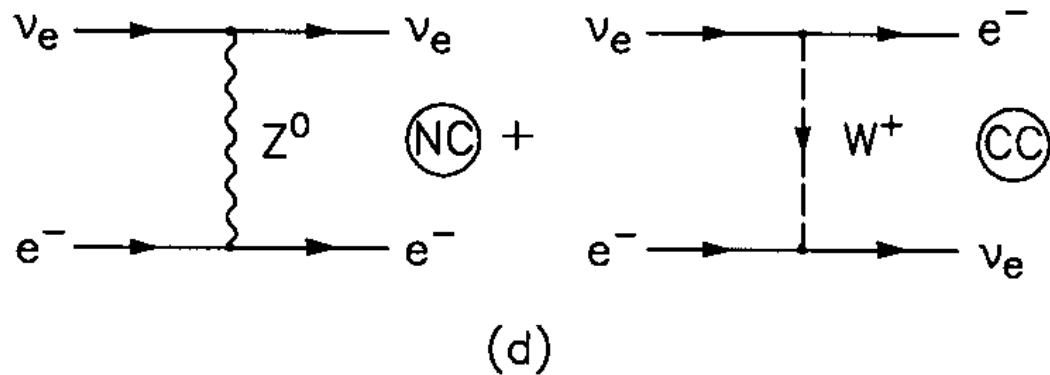
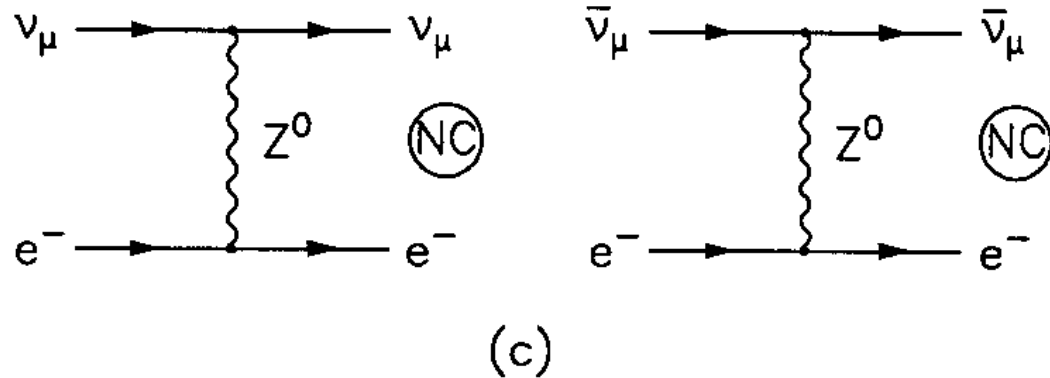
Neutrinos: $P(\nu_{e_l} \rightarrow \nu_{e_m}) = P_{CP} + P_{CP}$
Antineutrinos: $P(\bar{\nu}_{e_l} \rightarrow \bar{\nu}_{e_m}) = P_{CP} - P_{CP}$

\Rightarrow e.g. CP Asymmetries:

$$a^{\text{CP}} := \frac{P(\nu_{e_l} \rightarrow \nu_{e_m}) - P(\bar{\nu}_{e_l} \rightarrow \bar{\nu}_{e_m})}{P(\nu_{e_l} \rightarrow \nu_{e_m}) + P(\bar{\nu}_{e_l} \rightarrow \bar{\nu}_{e_m})} = \frac{P_{CP}}{P_{CP}}$$

Matter Effects and MSW Resonance

Mikheyev-Smirnov-Wolfenstein: **coherent forward scattering**



$\mathcal{L}_{NC} = \text{flavour universal}$

$\mathcal{L}_{CC} = \sqrt{2}G_F n_e \Leftrightarrow \text{only } \nu_e$

MSW-resonance energy (Δm_{31}^2)

Earth: $E_{\text{res}} \simeq 10 \text{ GeV}$

dominated by average density

$$\rho = \rho_{\text{average}} + \delta\rho$$

Baseline & MSW Matter Effects

Beams in Earth Matter:

⇒ electron density profile
as function of radius **Stacy**
density errors **Geller & Hara**

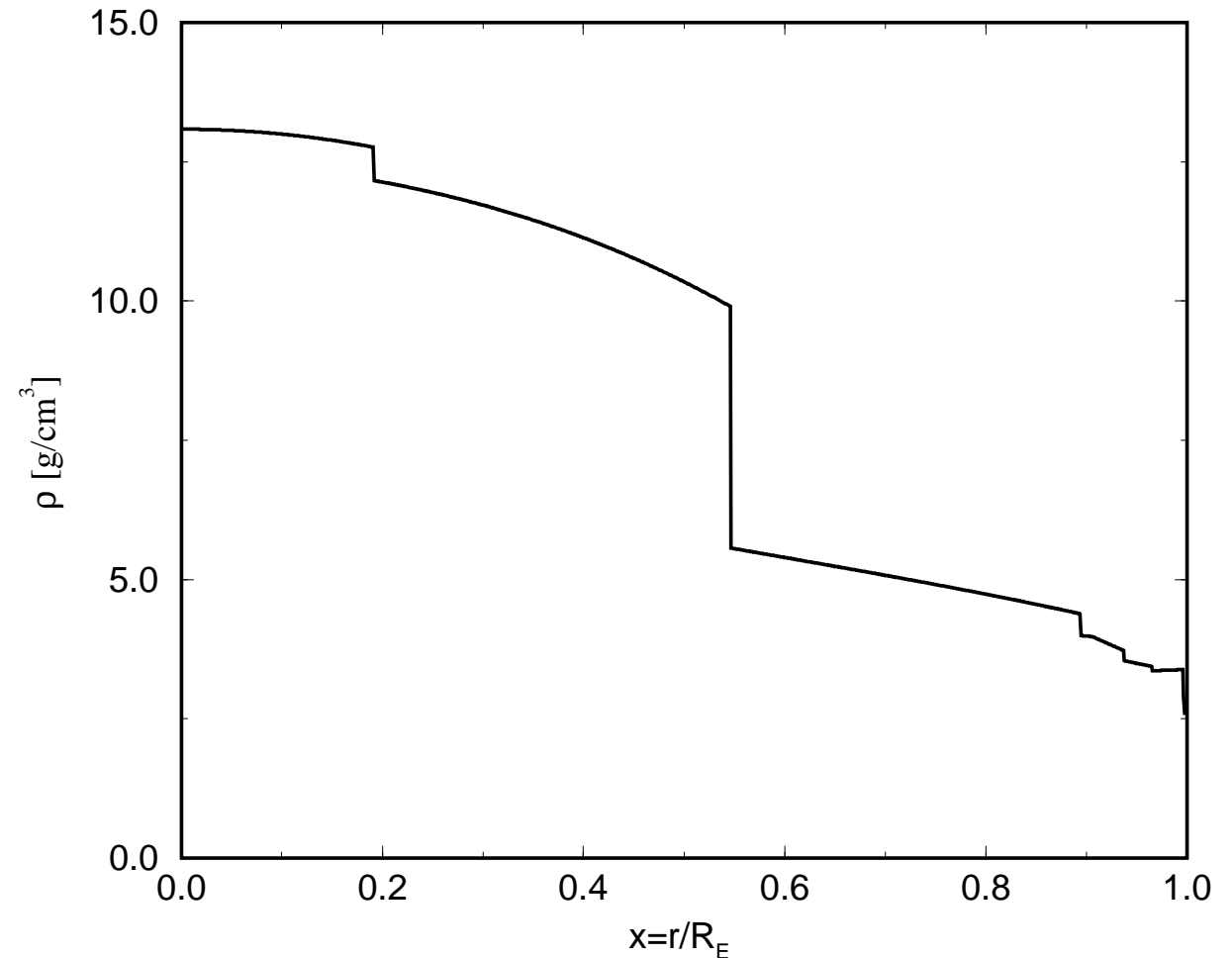
Large L ⇒ steep angles

$L = 2800 \text{ km} \Leftrightarrow 13^\circ$

$L = 7300 \text{ km} \Leftrightarrow 35^\circ$

$L = 12750 \text{ km} \Leftrightarrow 90^\circ$

$L \leq \mathcal{O}(10\,000 \text{ km}) \Leftrightarrow$ **mantle**



- $E_{\text{resonance}} \simeq 10 - 15 \text{ GeV}$, **matter effects grow with distance L**
- **Average density** profile uncertainties **decrease with L ⇒ $\simeq 5\%$ error**

$\Delta m_{12}^2 \simeq 0$, $Y = e^-/\text{nucleon}$ $\rho = \text{matter density}$ $m_n = \text{nucleon mass}$

In Flavour Basis:

$$H_0 + \delta \mathbf{H}_{CC} + \delta \mathbf{H}_{NC} = \frac{1}{2E} \mathbf{U} \begin{pmatrix} m_1^2 & 0 & 0 \\ 0 & m_2^2 & 0 \\ 0 & 0 & m_3^2 \end{pmatrix} \mathbf{U}^T + \frac{1}{2E} \begin{pmatrix} \mathbf{A} + \mathbf{A}' & 0 & 0 \\ 0 & \mathbf{A}' & 0 \\ 0 & 0 & \mathbf{A}' \end{pmatrix}$$

$$\mathbf{A} = \pm \frac{2\sqrt{2}G_F Y \rho E}{m_n},$$

“+” for $\nu \oplus \text{matter}$ and $\bar{\nu} \oplus \overline{\text{matter}}$

$\mathbf{U} = R_{23}R_{13}R_{12} \rightarrow \mathbf{U}' \Rightarrow$ **1-3 subspace parameter mapping**

$$\begin{aligned} \sin^2 2\theta'_{13} &= \frac{\sin^2 2\theta_{13}}{C_{\pm}^2} \\ \Delta m_{31,m}^2 &= \Delta m^2 C_{\pm} \\ \Delta m_{32,m}^2 &= \frac{\Delta m^2 (C_{\pm} + 1) + A}{2} \\ \Delta m_{21,m}^2 &= \frac{\Delta m^2 (C_{\pm} - 1) - A}{2} \end{aligned}$$

$$C_{\pm}^2 = \left(\frac{A}{\Delta m^2} - \cos 2\theta \right)^2 + \sin^2 2\theta$$

\Rightarrow **Different mappings for neutrinos and antineutrinos**

⇒ **All CP-violating effects are proportional to:**

1) $8J_{\text{CP}} = \cos \theta_{13} \sin(2\theta_{13}) \sin(2\theta_{12}) \sin(2\theta_{23}) \sin \delta$

2) $\sin \Delta_{12} \sin \Delta_{23} \sin \Delta_{31} \Rightarrow$ **3 different masses required**

P_{CP} is suppressed by:

- **small solar Δm_{sun}^2 :** $\sin \Delta_{12} \approx \Delta_{12} \ll 1$
- **small mixing angles:** $\sin^2 2\theta_{13} \ll 1$
- **in addition for SMA:** $\theta_{\text{sun}} \simeq \theta_{12}$

P_{CP} is not suppressed by $\Delta_{12} = \frac{\Delta m_{12}^2 L}{4E}$

Sizable CP effects:

3 Neutrinos \Rightarrow LMA-MSW (most likely)
4 Neutrinos \Rightarrow always

Qualitative Analytic Understanding

- full numerical simulation
- $\Delta = \Delta m_{31}^2 L/4E$
- qualitative understanding \Rightarrow expand in $\alpha = \Delta m_{21}^2/\Delta m_{31}^2$ and $\sin^2 2\theta_{13}$
- matter effects $\hat{A} = A/\Delta m_{31}^2 = 2VE/\Delta m_{31}^2$; $V = \sqrt{2}G_F n_e$

$$P(\nu_\mu \rightarrow \nu_\mu) \approx 1 - \cos^2 \theta_{13} \sin^2 2\theta_{23} \sin^2 \Delta + 2 \alpha \cos^2 \theta_{13} \cos^2 \theta_{12} \sin^2 2\theta_{23} \Delta \cos \Delta$$

$$\begin{aligned}
 P(\nu_e \rightarrow \nu_\mu) &\approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2((1-\hat{A})\Delta)}{(1-\hat{A})^2} \\
 &\pm \sin \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\
 &+ \cos \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\
 &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

Freund

Cervera & Donini & Gavela & Gomez Cadenas & Hernandez & Mena & Rigolin

Freund & Huber & ML

The Importance of θ_{13}

CHOOZ \Rightarrow θ_{13} is small: $\sin^2 2\theta_{13} < 0.1$

All effects in the $\nu_e \rightarrow \nu_\mu$ -transition depend crucially on θ_{13} :

- the total transition rate
- matter effects
- the effects due to the sign of Δm_{31}^2
- CP violating effects

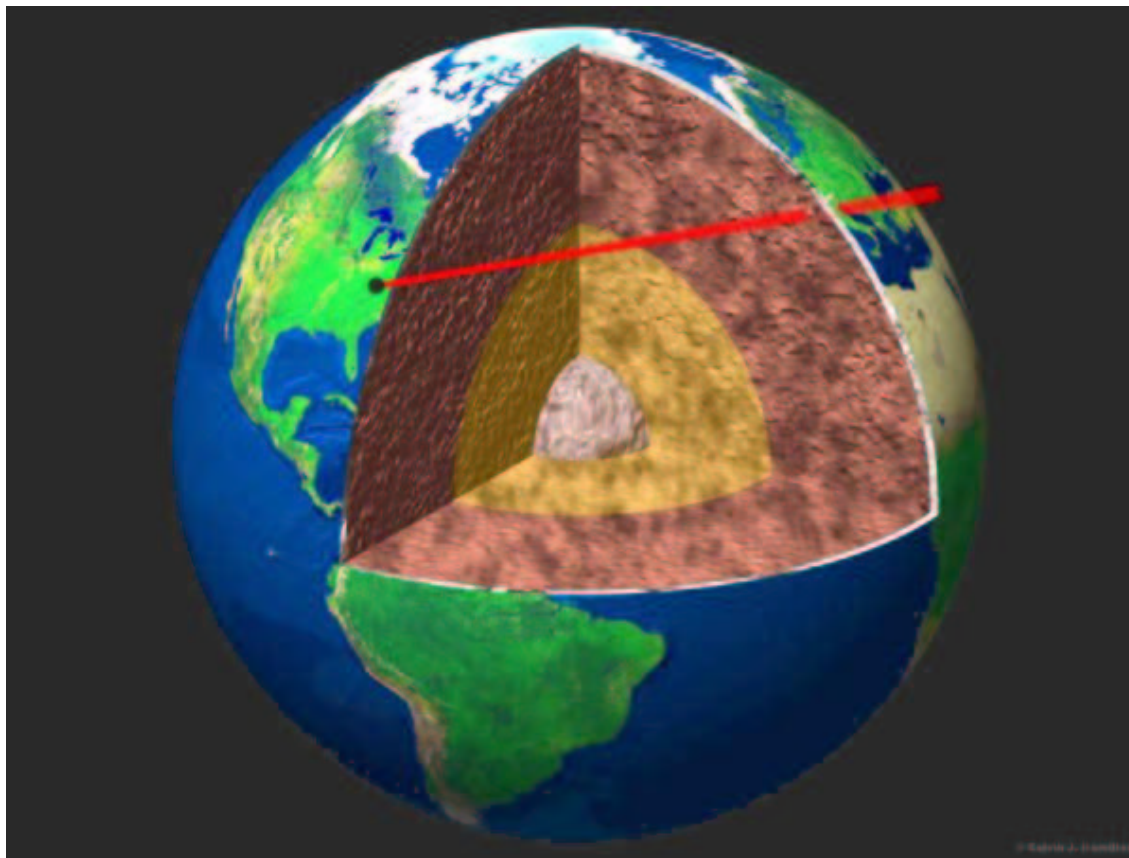
The size of θ_{13} determines if these effects can be studied

Errors, Degeneracies and Correlations

Look at expansion in powers of $\alpha = \Delta m_{21}^2 / \Delta m_{31}^2$ and $\sin \theta_{13}$; $\Delta = \Delta m_{31}^2 L / 4E$; $V = 0$

$$\begin{aligned}
 P(\nu_e \rightarrow \nu_\mu) &\approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2(\Delta) \\
 &\pm \sin \delta_{CP} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \sin^3(\Delta) \\
 &+ \cos \delta_{CP} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\Delta) \sin^2(\Delta) \\
 &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \sin^2(\Delta)
 \end{aligned}$$

- depends on solar parameters only via the **product** $\Delta m_{21}^2 \cdot \sin 2\theta_{12}$
- most interesting: $\sin^2 2\theta_{13} \simeq \alpha \cdot \sin^2 2\theta_{12} = \sin^2 2\theta_{12} \Delta m_{21}^2 / \Delta m_{31}^2$
- **degeneracies**: $(\delta_{CP} - \theta_{13})$, $(\delta_{CP} - \text{sign}(\Delta m_{31}^2))$, $(\theta_{23} - \frac{\pi}{2} - \theta_{23})$
- $\alpha^2 = (\Delta m_{21}^2)^2 / (\Delta m_{31}^2)^2$ term dominates **for tiny** $\sin^2 2\theta_{13} \Leftrightarrow$ **error of** Δm_{21}^2



Precision!

Source



Oscillation



Detector

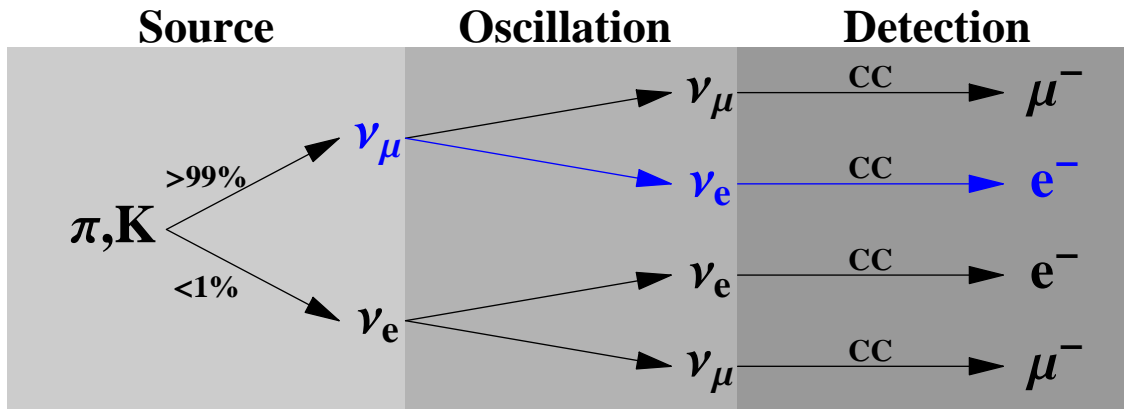
- neutrino energy E
- flux and spectrum
- flavour composition
- contamination
- symmetric $\nu/\bar{\nu}$ operation

- oscillation channels
- realistic baselines
- MSW matter profile

- effective mass, material
- threshold, resolution
- particle ID (flavour, charge, event reconstruction, ...)
- backgrounds
- x-sections (at low E)

Sources, Oscillation and Detection

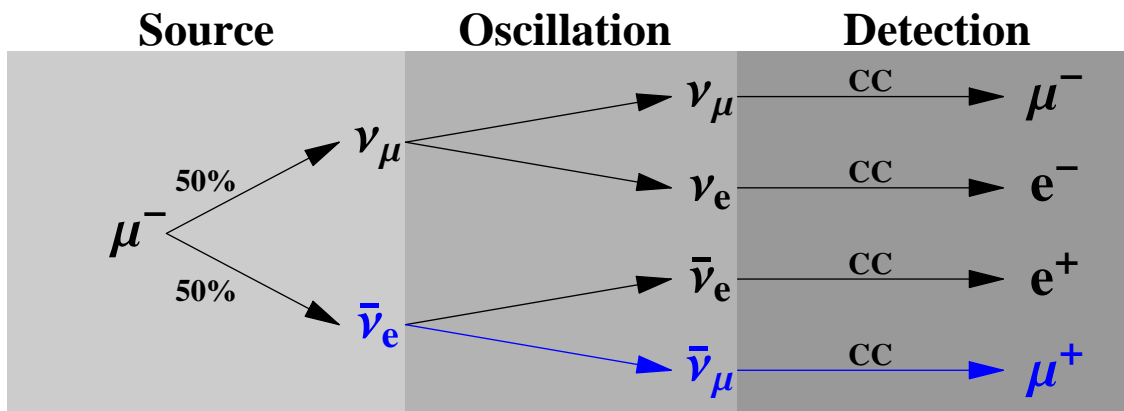
A) Conventional ν -Beams from Beam Dumps \Rightarrow Superbeams



$\nu_\mu \rightarrow \nu_e$ **oscillation** most interesting
 ν_e **contamination** \Leftrightarrow **off-axis**
good electron detection efficiency
good NC background rejection
near detector
 $\bar{\nu}$ -beam \simeq **different experiment**

B) Neutrino Factories

Geer



$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$ **oscillation** most interesting
excellent beam properties
very good charge ID required
good NC background rejection
 μ^+ mode **very symmetric**

C) Radioactive Beams ? **Pure ν_e or $\bar{\nu}_e$ beam** from β decay, $\gamma \simeq 100$ Zuchelli

The Potential of Future LBL Setups

Physis input:

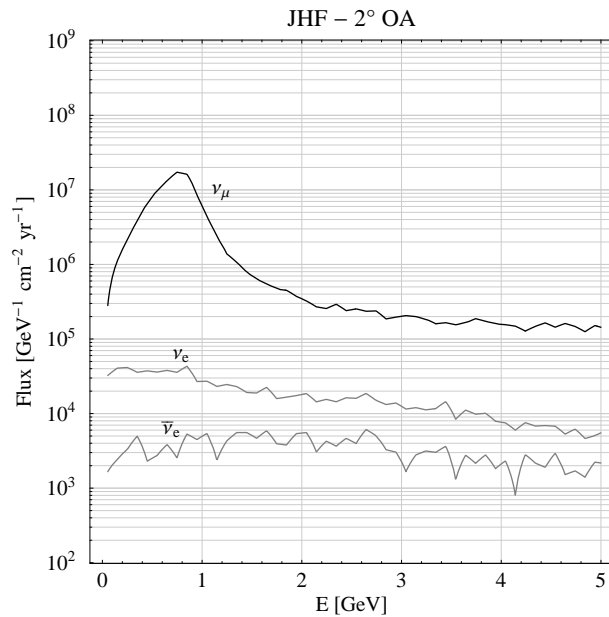
- **full numerical 3- ν simulation** and analytic understanding
- **matter density profile** and errors taken into account
- **error correlations included** (no fixed unknown parameters)

Analysis:

- **MC for all possible input parameters**
- **combined fit of appearance and disappearance channels for both polarities**
- **fit spectral information and/or total rates**
- **adequate statistical methods** (small event rates/bin):
 - ⇒ **Poissonian statistics**
 - ⇒ **parameterization of systematical uncertainties**
 - ⇒ **integrating out the nuisance parameters**
 - ⇒ **projection on the parameter of interest**
 - ⇒ **external information from KamLand (=LMA) and geophysics**
 - ⇒ **6 remaining parameters**
- ⇒ **Extract parameters and (correlated) errors** \Leftrightarrow **sensitivity**

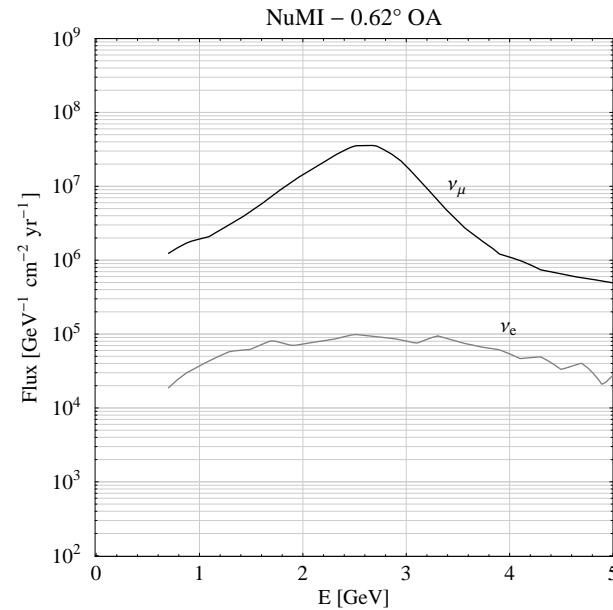
The Sources

JHF



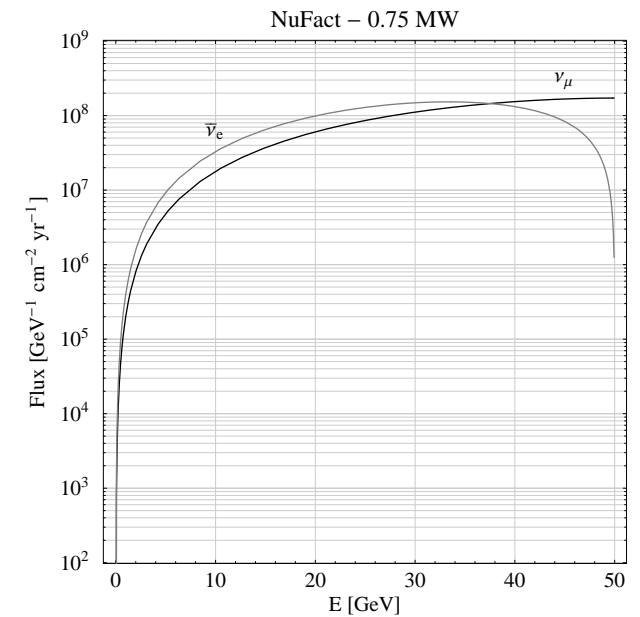
- **mean energy** 0.51 GeV
- **peak intensity**
 $1.7 \cdot 10^7 \text{ GeV}^{-1} \text{ cm}^{-2} \text{ yr}^{-1}$
at 0.78 GeV
- **ν_μ/ν_e -ratio at peak** 0.2%

NuMI off-axis



- **mean energy** 2.78 GeV
- **peak intensity**
 $3.6 \cdot 10^7 \text{ GeV}^{-1} \text{ cm}^{-2} \text{ yr}^{-1}$
at 2.18 GeV
- **ν_μ/ν_e -ratio at peak** 0.2%

NuFact



- **mean energy** 30 GeV
- **peak intensity**
 $1.5 \cdot 10^8 \text{ GeV}^{-1} \text{ cm}^{-2} \text{ yr}^{-1}$
at 33.33 GeV
- **ν_μ/ν_e -ratio at peak** 83%

Uncertainties in flux and ν_e -background

Itow et al.

Para, Szleper

Uncertainties in flux

Geer

Detectors

	water Cherenkov = SK (HK)	low-Z calorimeter	magnetized iron calorimeter
fiducial mass	22.5 kt (1 000 kt)	20 kt	10 kt (50 kt)
energy range	0.4 – 1.2 GeV	1 – 5 GeV	4 – 50 GeV
energy resolution	5%	10%	20%
signal efficiency	0.5	0.5	0.45
NC rejection	0.01	0.001	$< 10^{-5}$
CID	–	–	$< 10^{-5}$
background uncertainty	5%	5%	5%

- **threshold effects** for the magnetized iron calorimeter
linear rise of the **efficiency** between 4 GeV and 20 GeV
- **liquid Argon TPC ?**

5 Considered Scenarios

	JHF-SK	NuMI	NuFact-I
detector	water cherenkov	low-Z calorimeter	10kt magnetized iron calorimeter
baseline	295 km	735 km	3 000 km
matter density	2.8 g cm^{-3}	2.8 g cm^{-3}	3.5 g cm^{-3}
L/E_{peak}	378 km GeV^{-1}	337 km GeV^{-1}	90 km GeV^{-1}

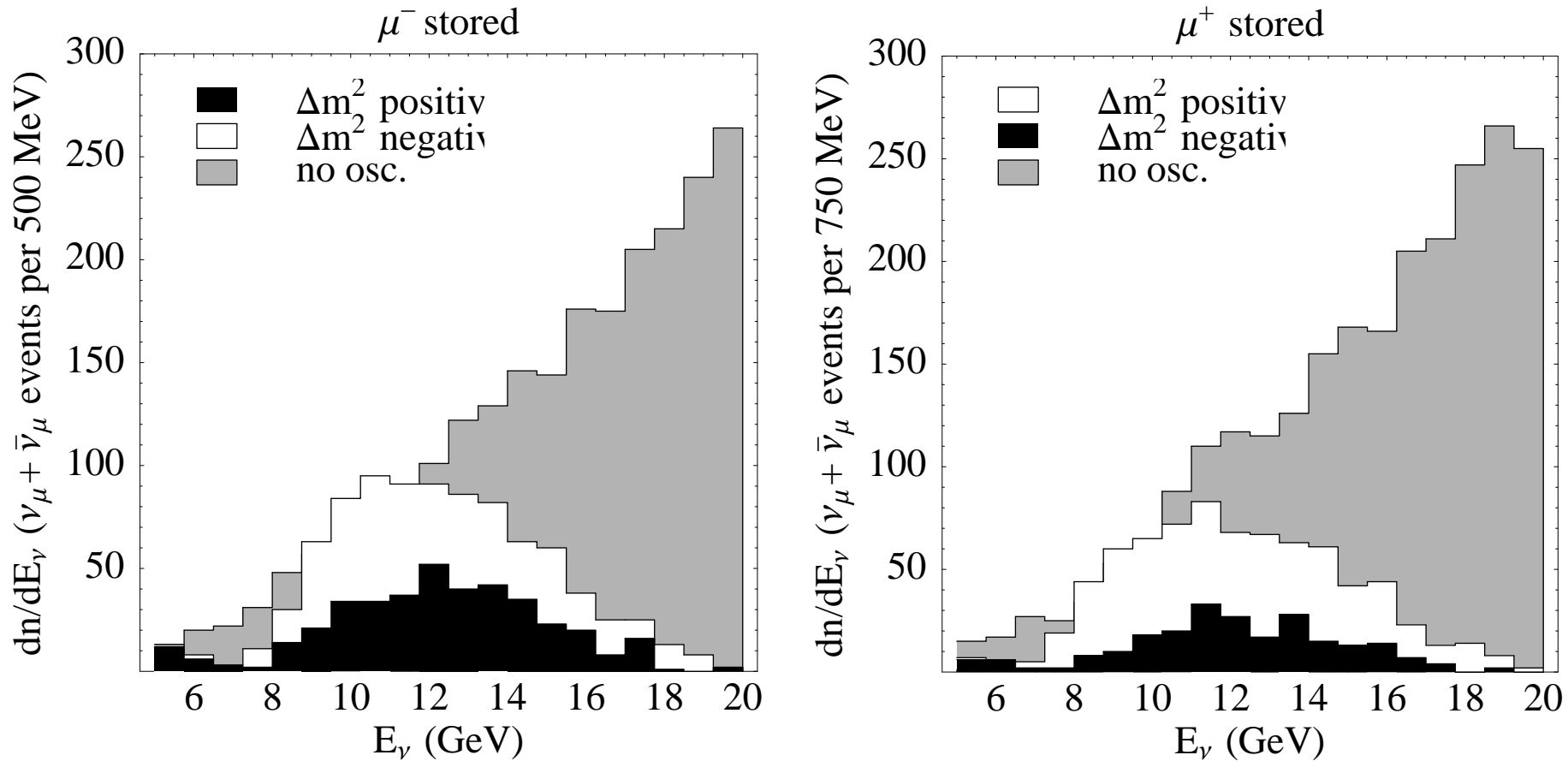
- **JHF-HK** $\Rightarrow \simeq 95$ times more luminosity than JHF-SK plus anti-neutrino running
- **NuFact-II** $\Rightarrow \simeq 42$ times more luminosity than NuFact-I

Itow et al., A. Para, A. Blondel et al.

	JHF-SK	NuMI	JHF-HK	NuFact-I	NuFact-II
signal	139.0	387.5	13 180.0	1 522.8	64 932.6
background	23.3	53.3	2 204.6	4.2	180.3
S/N	6	6	6	360	360

Disappearance Channels: Qualitative Picture

Large event rates \Rightarrow spectral information \Rightarrow NuFact example:



$\Rightarrow \Delta m_{31}^2, \theta_{23}$ and for not too small θ_{13} also $sign(\Delta m_{31}^2)$

Cervera et al., Barger & Geer & Raja & Whisnant, DeRujula et al., Donini et al., Bueno & Campanelli & Rubbia, Albright et al., Yasuda, Minakata, Shrock et al., Barenboim & De Gouvea & Sziper & Velasco, Freund & ML & Petcov & Romanino, Freund & Huber & ML,

Appearance Channels: Qualitative Picture

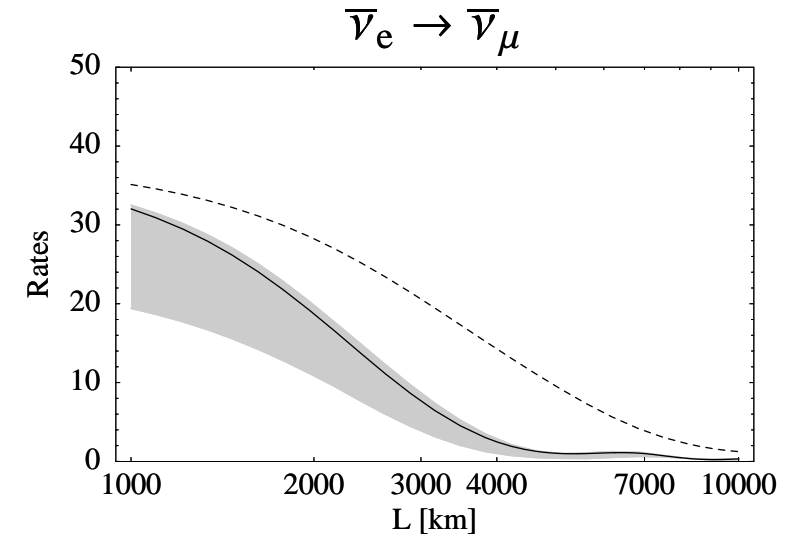
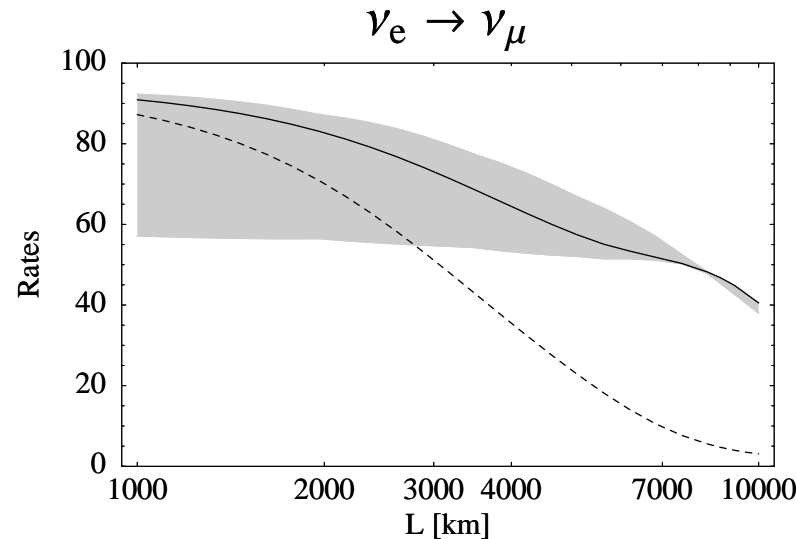
Small event rates \Rightarrow total rates & little spectral information

NuFact:

$E_\mu = 50$ GeV

$\sin^2 2\theta_{13} = 0.01$

LMA



solid lines=matter ($\delta = 0$), dashed lines=vacuum ($\delta = 0$), grey band=all CP phases

Comparable matter and CP effects \Rightarrow separation
CP-phase \Rightarrow short(er) baseline
matter effects, sign of $\delta m^2 \Rightarrow$ large L

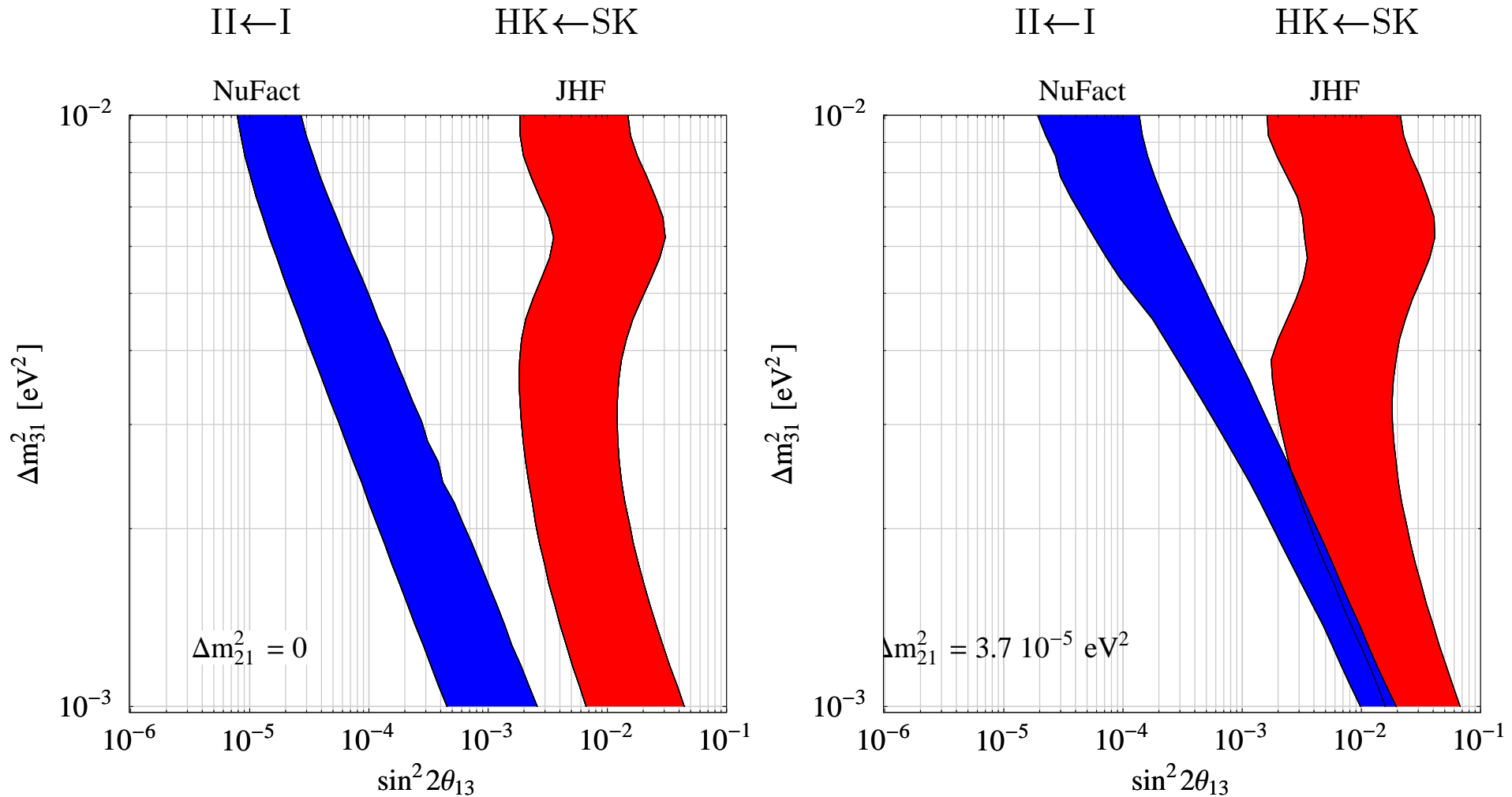
For all solar solutions: θ_{13} and maybe $sign(\Delta m_{31}^2)$

For LMA without external input: $\simeq \theta_{12}$, $\simeq \Delta m_{21}^2$

With external input for θ_{12} and Δm_{21}^2 : $\Rightarrow \delta_{CP}$

Sub-Leading Parameters: $\sin^2 2\theta_{13}$ Sensitivity

All parameter correlations must be taken into account



- sizable Δm_{31}^2 and Δm_{21}^2 dependence of sensitivity limits
- up to an order of magnitude difference
- effect strongest for short baselines

Impact of Systematical Errors

Signal

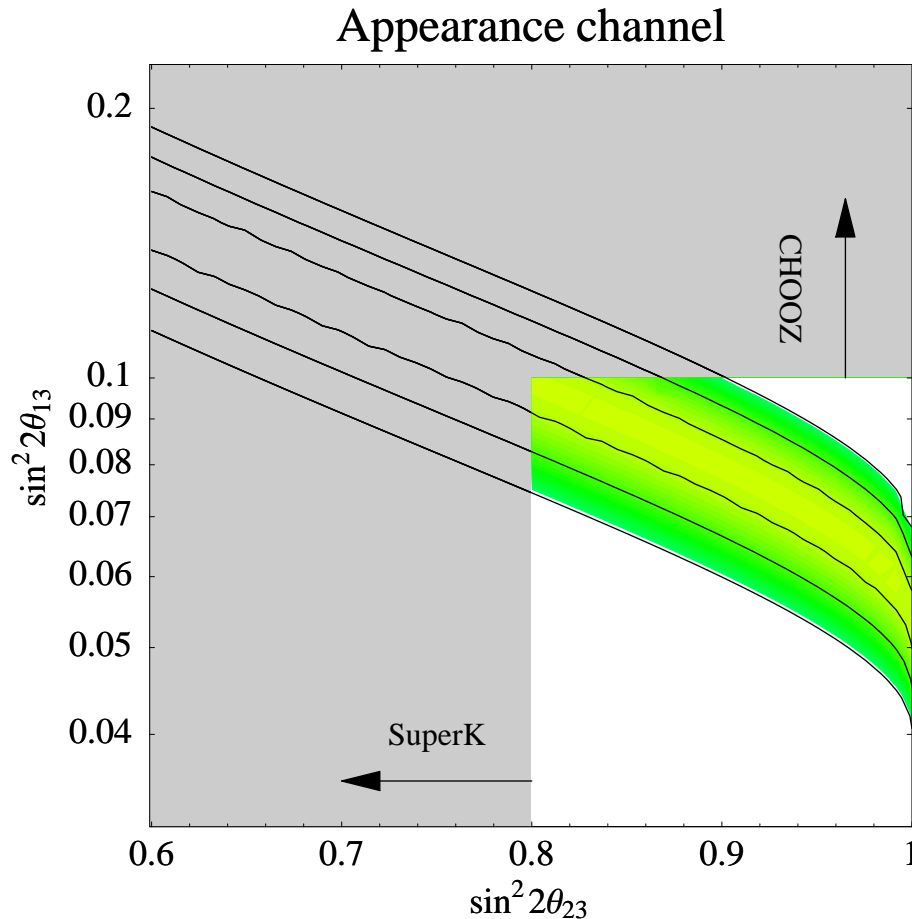
- use the **information in the energy spectrum**
- overall **normalization error** of 5% is acceptable
- **energy calibration error** of 5% is acceptable

Background

- **backgrounds limit sensitivity**
- **energy spectrum** helps to reduce the impact of the background
- background uncertainty becomes important when $\sigma_{b_n} \geq 1/\sqrt{b}$
for $\sigma_{b_n} = 5\% \Leftrightarrow b = 400$

Background uncertainties may dominate high statistics experiments

Impact of Correlations



⇐ **Example**

- multi-parameter problem
- highly non-linear
- complex topology

⇒ **use all available information:**

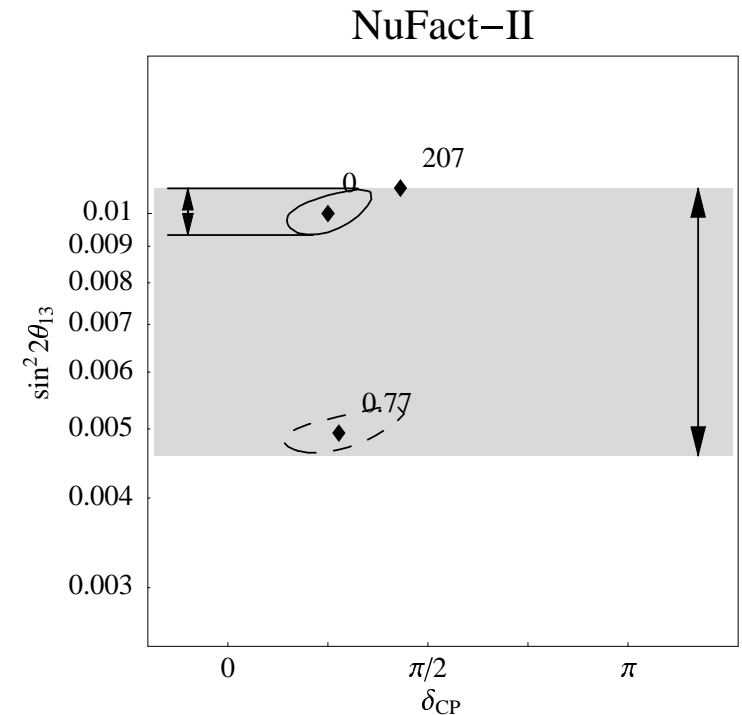
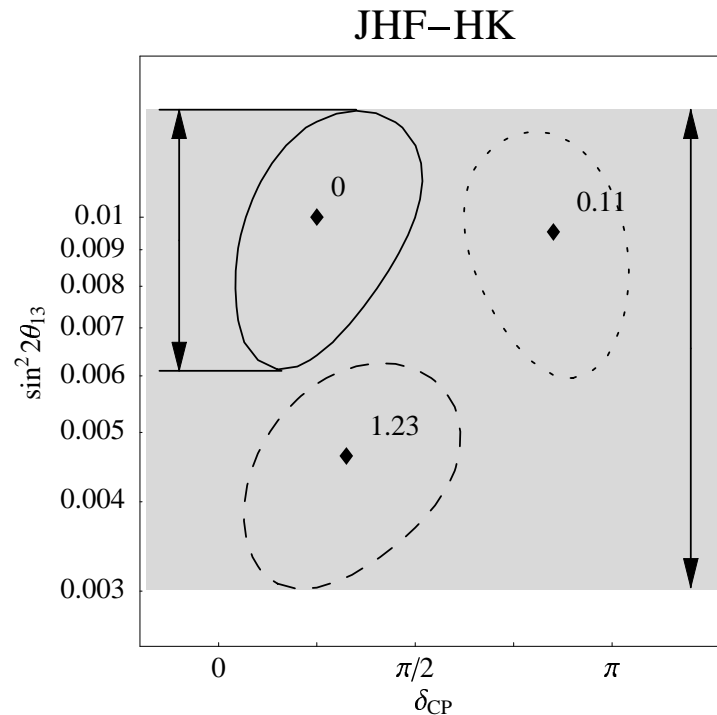
- appearance & disappearance channel
- energy information
- external input

Impact of Degeneracies

Complex parameter dependence in appearance probabilities \Rightarrow **multiple solutions**

3 degeneracies:

- $\delta_{\text{CP}} - \theta_{13}$
- $\delta_{\text{CP}} - \text{sgn} \Delta m_{31}^2$
- $\theta_{23} - \pi/2 - \theta_{23}$



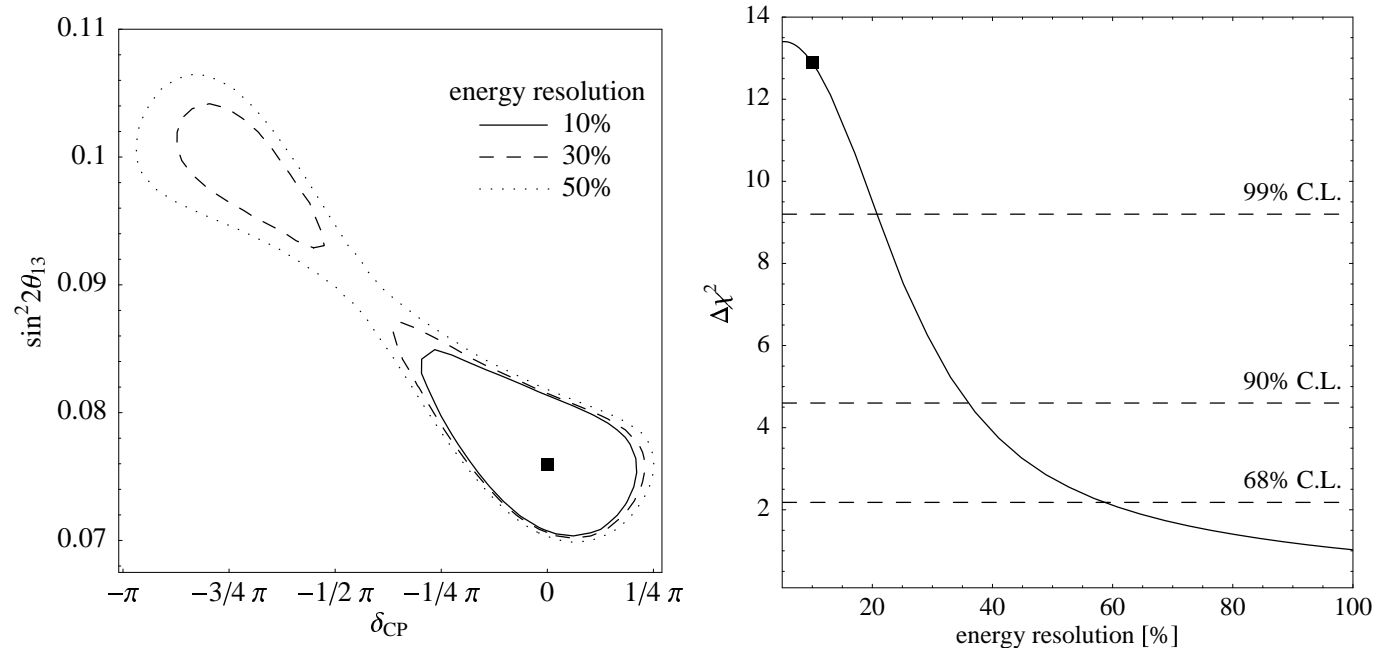
Full analysis:

- $\delta_{\text{CP}} - \theta_{13}$ is transformed to a correlation \Leftrightarrow neutrino beam only
- $\delta_{\text{CP}} - \text{sgn} \Delta m_{31}^2$ remains
- $\theta_{23} - \pi/2 - \theta_{23}$ remains
- **combined or separate regions**

Cervera et al., Barger et al., Burguett-Castell et al., Minakata et al., Huber, ML, Winter

Medium Baseline & Good Resolution

Correlation of δ_{CP} with θ_{13}

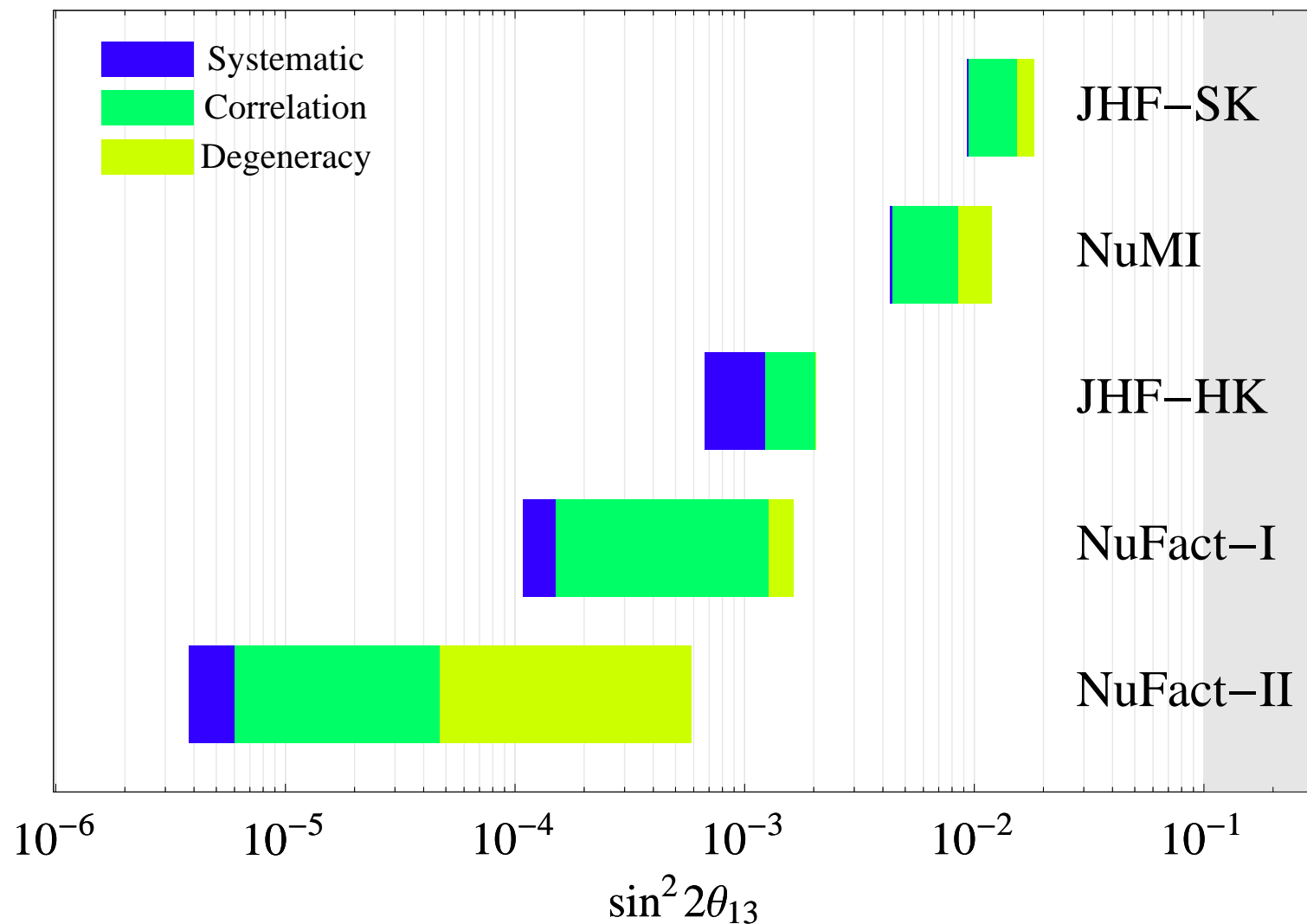


$$E = 50 \text{ GeV}, L = 3000 \text{ km}, \Delta m_{21}^2 = 10^{-4} \text{ eV}^2$$

- two degenerate solutions
- depends strongly on energy resolution
- degeneracy lifted for energy resolution better than 25%

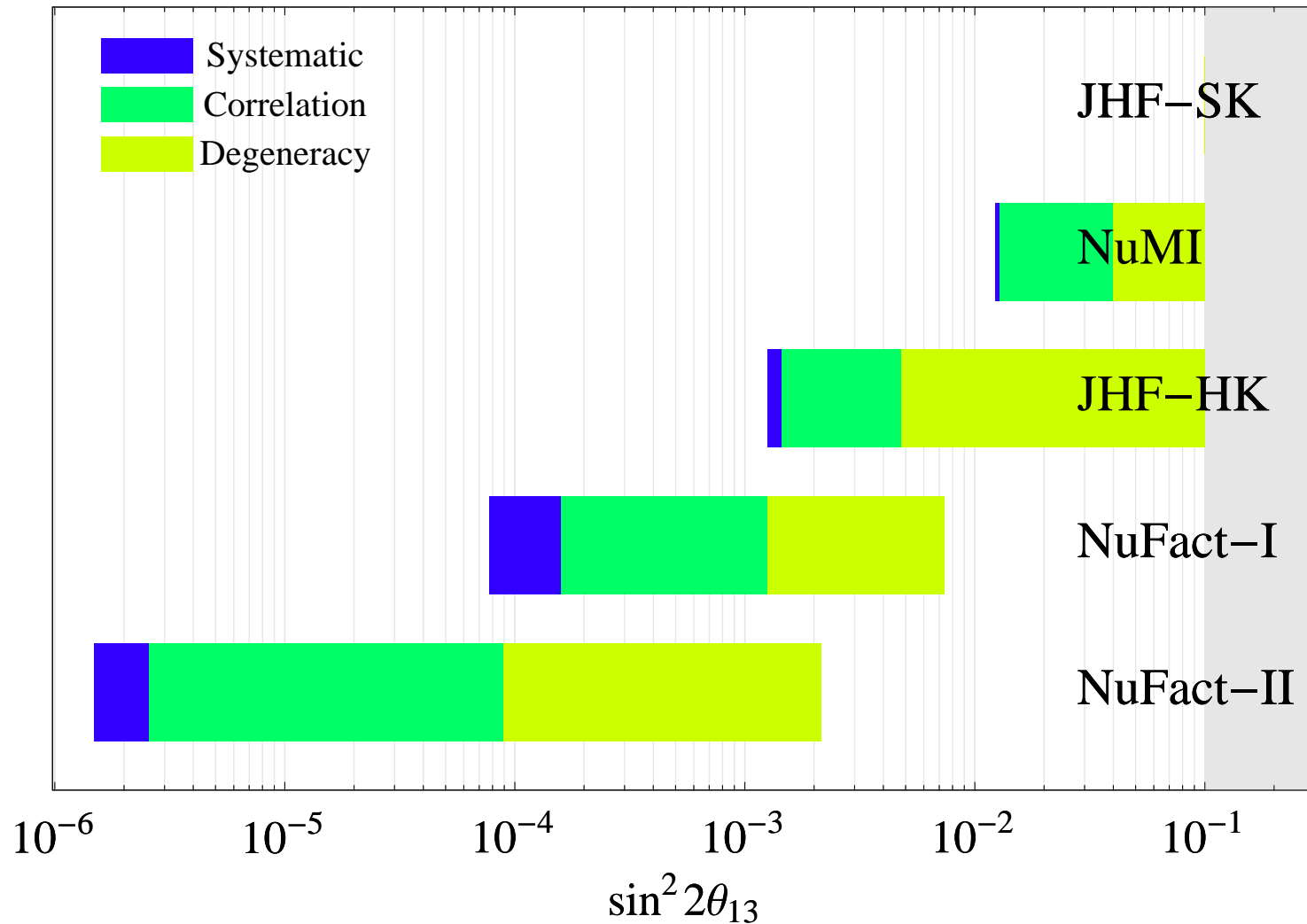
degeneracy can be resolved with spectral information \Leftrightarrow resolution & flux

Sensitivity to $\sin^2 2\theta_{13}$



- Different sensitivity reductions by systematics
- Correlations & degeneracies lead to severe limitations
- Improvements by combining experiments

Sensitivity to the sign of Δm_{31}^2

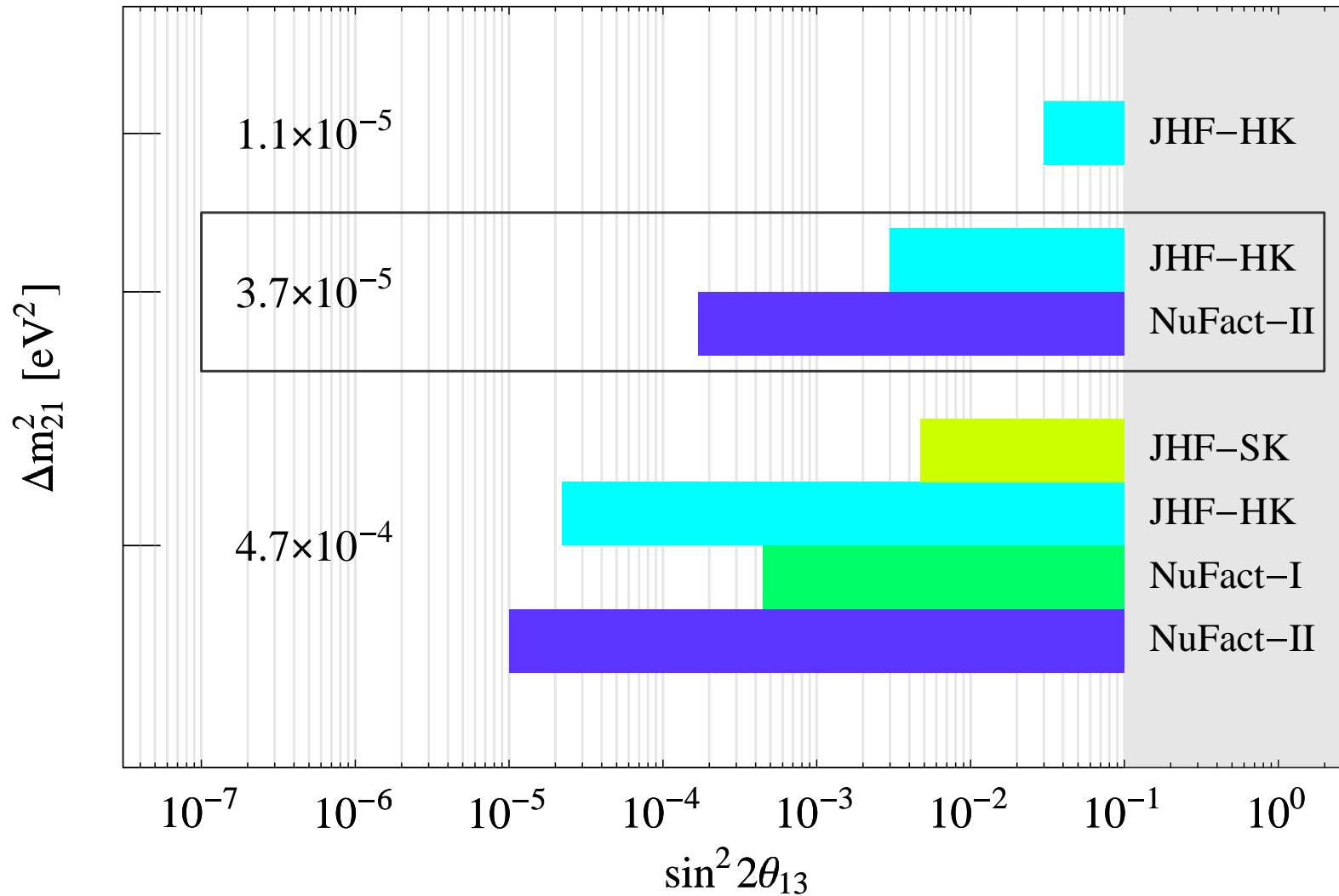


- $sign(\Delta m_{31}^2)$ very hard to determine with superbeams
 - **degeneracies with δ_{CP}** are the main problem
- ⇒ **combine experiments!**

Huber, ML, Winter, hep-ph/0204352

Measurements of CP-violation

Sensitivity to CP-Violation at $\delta_{CP} = +\pi/2$



• **CP violation** with high luminosity superbeams **feasible**

• **sensitivity is δ_{CP} dependent**

Huber, ML, Winter, hep-ph/0204352

T-Violating Effects

E. Akhmedov, P. Huber, M. Lindner, and T. Ohlsson, NPB 608, 394 (2001), hep-ph/0105029.

CP, T and CPT Differences:

$$\Delta P_{\alpha\beta}^{CP} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

$$\Delta P_{\alpha\beta}^T \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\nu_\beta \rightarrow \nu_\alpha)$$

$$\Delta P_{\alpha\beta}^{CPT} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$$

CPT is conserved in local QFT \Rightarrow

$$P(\nu_\alpha \rightarrow \nu_\beta) \equiv P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$$

$$\Delta P_{\alpha\beta}^{CP} + \Delta P_{\bar{\alpha}\bar{\beta}}^T = \Delta P_{\alpha\beta}^{CPT} = 0$$

CPT & vacuum \Rightarrow correlated CP and T violation from NMS matrix

Oscillations in Matter: e^- and nucleons, i.e. particles, no antiparticles \Rightarrow

CP and CPT are violated by matter

\Rightarrow Non-trivial interplay of **fundamental** and **matter-induced** T violation

Experimental Problem:

- No direct test of T-violation: **Direction of time cannot be changed**
- Alternative: ν -osc. forward in time with **initial and final flavors interchanged**
 \Rightarrow Interchange of source and detector \Leftrightarrow **asymmetric density profile**

T violation in neutrino oscillations:

fundamental (intrinsic) T violation \Leftrightarrow non-vanishing δ_{CP}

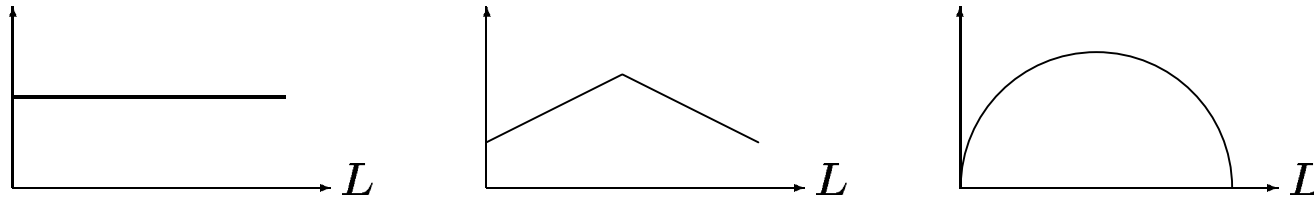
matter-induced T violation \Leftrightarrow asymmetric profiles

$$\Delta P_{\alpha\beta}^T \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\nu_\beta \rightarrow \nu_\alpha)$$

In Vacuum: CPT \Rightarrow T violation \Leftrightarrow CP violation

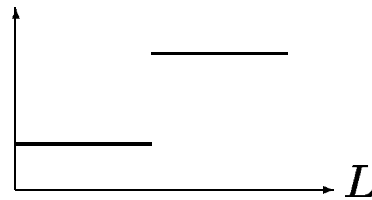
Matter: 2 ν -flavors \Rightarrow NO T-violating effects: $P_{e\mu} = P_{\mu e} \Rightarrow \Delta P_{e\mu} = 0$

Matter: 3 ν -flavors, symmetric matter density profiles



\Rightarrow If $\delta_{CP} = 0$, then $\Delta P_{\alpha\beta}^T = 0$

Matter: 3 ν -flavors, asymmetric matter density profiles



\Rightarrow General Case

T-odd probability difference: Arbitrary density profile, small θ_{13} and δ/Δ

$$\begin{aligned}\Delta P_{e\mu}^T &\simeq -2s_{23}c_{23} \text{Im} [\beta^*(A - C^*)] \\ &\simeq -2s_{13}s_{23}c_{23} (\Delta - s_{12}^2\delta) \text{Im} [e^{-i\delta_{CP}}\beta^*(A_a - C_a^*)]\end{aligned}$$

where

$$s_{ij} \equiv \sin \theta_{ij}, \quad c_{ij} \equiv \cos \theta_{ij};^1, \quad \delta \equiv \frac{\Delta m_{21}^2}{2E_\nu}, \quad \Delta \equiv \frac{\Delta m_{31}^2}{2E_\nu};$$

$$A_a \equiv \alpha \int_{t_0}^t \alpha^* f dt' + \beta \int_{t_0}^t \beta^* f dt', \quad C_a \equiv f \int_{t_0}^t \alpha f^* dt'.$$

$\alpha = \alpha(t, t_0)$, $\beta = \beta(t, t_0)$ given by solutions of 2-flavor problem in the (1,2)-sector and $f = f(t, t_0) = \exp \left\{ -i \int_{t_0}^t (\Delta - \frac{1}{2}[\delta + V(t')]) dt \right\}$

In addition: $\Delta P_{e\mu}^T = \Delta P_{\mu\tau}^T = \Delta P_{\tau e}^T$

$\Delta P_{e\mu}^T$ has been explicitly calculated for

- 1) matter consisting of **two layers of constant density** and
- 2) **arbitrary matter density profile** in the adiabatic approximation.

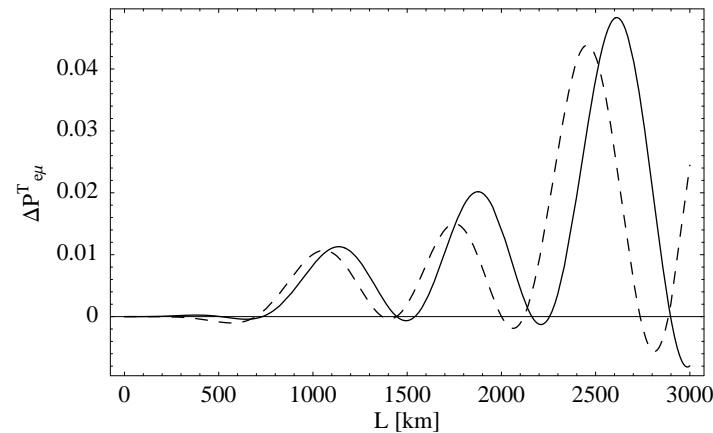
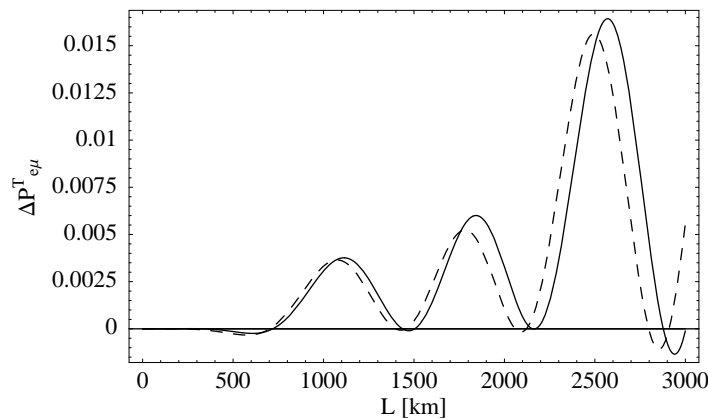
¹ $\theta_{12} = \theta_3, \theta_{13} = \theta_2, \theta_{23} = \theta_1$

Low energy regime ($\delta = \Delta m_{21}^2/(2E_\nu) \gtrsim V_{1,2}$), two layers:

$$\Delta P_{\alpha\beta}^T \simeq \underbrace{\cos \delta_{CP} \cdot 8 \cdot s_{12} c_{12} s_{13} s_{23} c_{23}}_{=J_{\text{eff}}} \frac{\sin(2\theta_1 - 2\theta_2)}{\sin 2\theta_{12}} \times \{s_1 s_2 [Y - \cos(\Delta_1 L_1 + \Delta_2 L_2)]\} \\ + \sin \delta_{CP} \cdot 4 s_{13} s_{23} c_{23} X_1 [Y - \cos(\Delta_1 L_1 + \Delta_2 L_2)]$$

$\cos \delta_{CP} \simeq$ matter-induced T violation, $\sin \delta_{CP} \simeq$ fundamental T violation

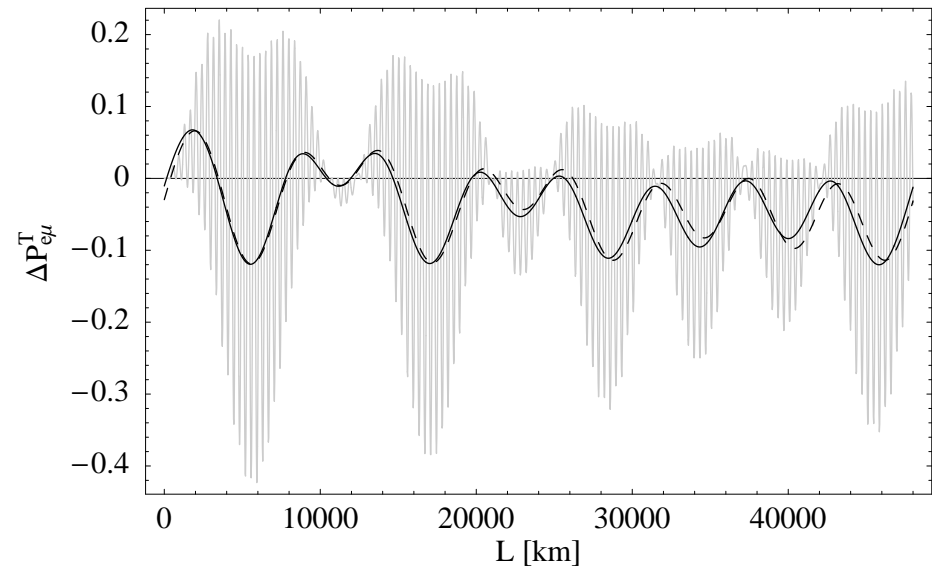
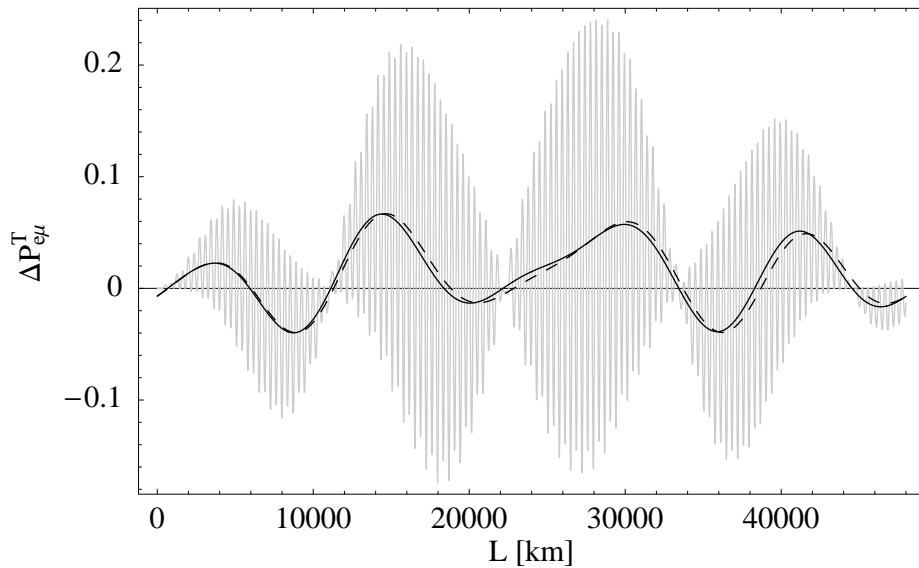
$\Delta P_{e\mu}^T(L)$:



$$L_1 = L_2 = L/2; \quad \rho_1 = 1 \text{ g/cm}^3, \quad \rho_2 = 3 \text{ g/cm}^3, \quad E_\nu = 1 \text{ GeV}$$

solid curve - analytic result , dashed curve - numerical result

$$\Delta P_{e\mu}^T = \Delta P_{e\mu}^T(L):$$



$$L_1 = L_2 = L/2; \quad \rho_1 = 0, \quad \rho_2 = 6.4 \text{ g/cm}^3$$

- grey curves - analytic results
- black solid & dashed curves - averaged over fast oscillations

$E_\nu = 0.5 \text{ GeV} \Rightarrow$ Large $\Delta m_{\text{atm}}^2 = \Delta m_{31}^2$ oscillations are **very fast!**

Left plot: Parameters as in **P.M. Fishbane and P. Kaus, PLB 506, 275 (2001)**

Right plot: Larger values of θ_{13} and Δm_{21}^2 .

NuFact on Earth $\Leftrightarrow \simeq$ **symmetric profiles** \Leftrightarrow **theoretical error in studies**

CPT-Violating Effects

S. Bilenky, M. Freund, M. Lindner, T. Ohlsson and W. Winter, Phys. Rev. D 65 (2002) 073024

CPT-violation



Physics beyond local QFT



“Planck Scale Physics”, “extra dimensions”, ...

- CPT violation was used to accommodate LSND result
- Here: Precision measurements \Rightarrow **interesting limits**

Sensitivity to small CPT-Violating Effects at NuFact

CPT-invariance $\Rightarrow P(\nu_\alpha \rightarrow \nu_{\alpha'}) = P(\bar{\nu}_{\alpha'} \rightarrow \bar{\nu}_\alpha)$

CPT-violation \Rightarrow Masses and mixings differ for neutrinos and antineutrinos.

Neutrinos: m_i, U **anti-neutrinos:** \bar{m}_i, \bar{U}

Consequence: $P(\nu_\alpha \rightarrow \nu_{\alpha'}) \neq P(\bar{\nu}_{\alpha'} \rightarrow \bar{\nu}_\alpha)$

Limit CPT-violation in the disappearance channels at NuFact

- high event rates
- no beam contamination
- small matter effects
- large oscillation effects

\Rightarrow Exclusion limits for tiny CPT-violating effects

Consider ν - and $\bar{\nu}$ channels as independent experiments:

$$\delta \equiv |\Delta m_{32}^2 - \Delta \bar{m}_{32}^2|, \quad \epsilon \equiv |\sin^2 2\theta_{23} - \sin^2 2\bar{\theta}_{23}|$$

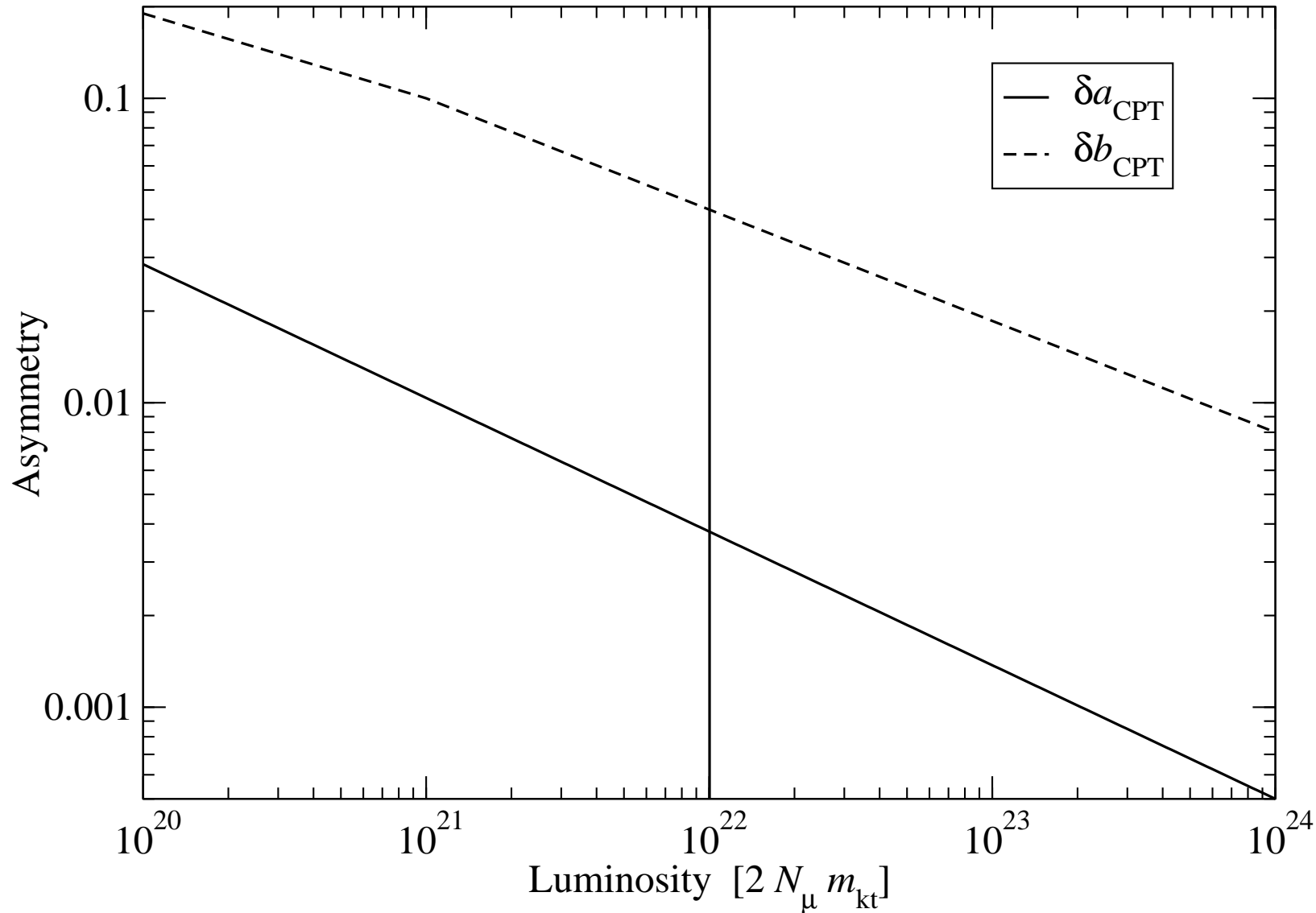
For $m_1 \ll \sqrt{\Delta m_{sun}^2}$ (hierarchical) and $|m_3 - \bar{m}_3| \ll (m_3)_{\text{average}}$

$$\delta \simeq 2 a_{\text{CPT}} \Delta m_{32}^2 \quad \text{with} \quad a_{\text{CPT}} \equiv \frac{|m_3 - \bar{m}_3|}{(m_3)_{\text{average}}}$$

$$\epsilon \simeq 2 b_{\text{CPT}} \sqrt{\sin^2 2\theta_{23}} \sqrt{1 - \sin^2 2\theta_{23}} \arcsin \sqrt{\sin^2 2\theta_{23}} \quad \text{with} \quad b_{\text{CPT}} \equiv \frac{|\theta_{23} - \bar{\theta}_{23}|}{(\theta_{23})_{\text{average}}}$$

- Sensitivity to possible CPT violation \Leftrightarrow accuracy of a_{CPT} and/or b_{CPT}
- Compare δ and ϵ with respective relative statistical error $\delta \Delta m_{32}^2$ or $\delta \theta_{23}$
- Sensitivities δa_{CPT} and δb_{CPT} for a_{CPT} and b_{CPT} are given by:

$$\delta a_{\text{CPT}} \sim \frac{\delta \Delta m_{32}^2}{2}$$
$$\delta b_{\text{CPT}} \sim \delta \theta_{23}$$



$E_\nu = 50\text{GeV}$, $L = 7000\text{km}$ (θ_{23}) / 3000km (Δm_{23}^2), 10^{20} muons/y, 5 years, 10kt

$$a_{CPT} \leq 3.8 \cdot 10^{-3}, \quad b_{CPT} \leq 4.3 \cdot 10^{-2} \quad \Leftrightarrow \quad |m_3 - \bar{m}_3| = 1.9 \cdot 10^{-4} \text{ eV}$$

Conclusions

CP, T and CPT violation in neutrino oscillations offer exciting perspectives!

CP violation:

- intrinsic CP violation from MNS matrix
- extra CP violation from matter effects
- both can be measured in future LBL experiments
- separation strategies ...

T violation:

- in vacuum correlated to CP violation
- complex interplay between fundamental and matter effects
- very nice theoretical features
- for LBL on earth small \Leftrightarrow theoretical precision limit

CPT violation:

- in local QFT not allowed
- matter induced CPT violation
- fundamental CPT violation \Leftrightarrow Planck scale physics